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An Algebraic Model for a Zinc/Bromine Flow Cell

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ABSTRACT

An algebraic model for a parallel plate flow cell is presented and used to predict various performance quantities, which are compared to those predicted by using previously published differential equation models. The results presented compare well with previous work. The model is based on the concept of using well-mixed zones and linear concentration and potential profiles for the diffusion layers and the separator. The Butler-Volmer equation is used for the electrochemical reactions, and the homogeneous reaction between bromine and bromide is included.

A zinc/bromine battery consists essentially of two electrodes, a porous separator, and a pair of storage tanks. The separator is placed between the electrodes in such a way that there is a thin gap between each of the electrodes and the separator. An electrolytic solution containing Zn\(^{2+}\) and Br\(^{-}\) (among other things) is forced through the gaps between the electrodes and the separator. After the solution from each gap exits the reactor, it enters the associated storage tank and is recycled later to the reactor.

During charge, the principle reaction at the anode is the oxidation of Br\(^{-}\) to Br\(_2\). At the cathode the principle reaction is the reduction of Zn\(^{2+}\) to Zn. An undesirable reaction at the cathode during charge is the reduction of Br\(_2\) to Br\(^{-}\). A homogeneous phase bulk reaction occurs between Br\(_3\) and Br\(^{-}\) to yield Br\(_2\). The two electrochemical reactions can be written in anodic form as

\[
\text{Br}^- \rightarrow \frac{1}{2} \text{Br}_2 + e^- \quad \text{Reaction 1}
\]

\[
\text{Zn} \rightarrow \frac{1}{2} \text{Zn}^{2+} + e^- \quad \text{Reaction 2}
\]

and the complexation reaction can be written as

\[
\text{Br}^- + \text{Br}_2 \rightarrow \text{Br}_3^- \quad \text{Reaction 3}
\]

The equilibrium constant for this homogeneous reaction is 17,000 (cm\(^3\)/mol) (1).

Previous authors have modeled the zinc/bromine flow cell (2-7). Lee and Selman (2, 3) developed a model to determine the current density distribution along the electrode surfaces. Lee (4) extended that work to include time variations of the current density distribution in an attempt to predict dendrite formation. Van Zee et al. (5) developed a simplified model to describe overall cell performance. Mader and White (6) also developed a model that can be used to predict the overall cell performance by introducing a “one-step” approximation. Evans and White (7) extended that work by including a porous electrode and the capability of predicting round-trip energy efficiencies. Evans and White (6) present a review of the work done on the zinc/bromine flow cell.

Two basic approaches have dominated the modeling of parallel plate cells. The simpler of these approaches is to assume that the cell is operated under diffusion-limited conditions (9). This method can accurately predict the reactor’s limiting current density and maximum conversions when the electrodes are operated at high overpotentials and the effects of ionic migration are negligible. The more general of these approaches is to write a series of governing equations that describe the reactor’s performance under all conditions. The equations written are a set of partial differential equations that represent conservation of mass and a single algebraic equation that represents the electroneutrality condition. After the equations are written, they are solved numerically subject to applicable boundary conditions. The numerical solution requires that the partial differential equations be approximated by suitable finite-difference equations. These equations are then solved simultaneously to determine the various concentrations and potentials. The fundamentals of this technique are presented by Newman (10) and others (11, 12).

The model presented herein is an algebraic simplification of the second approach. In the model development it is assumed that the concentration profiles within the reactor can be approximated by a series of constant and linear regions and that the potential profiles can be approximated by a series of linear regions shown in Fig. 1. The regions of constant concentration correspond with perfectly mixed regions in the bulk of the electrolyte. The regions where the concentration profiles are linear correspond with regions near the electrodes where both diffusion and migration are present. The potential profile is assumed to be linear through each region of the reactor. This algebraic approach to modeling a parallel plate flow cell results in a significant simplification over the more rigorous approach (11, 12). Specifically, the algebraic model results in far fewer dependent variables to be determined.

Model Development

The model presented here consists of a material balance for each component in every region and a charge balance for every region, as shown in Fig. 1. The major assumptions of this model are: (i) there is a diffusion region of known thickness near each electrode, (ii) there is a perfectly mixed region a distance away from each electrode, and (iii) the electrolyte composition undergoes a step change from the inlet to the outlet composition within the reactor. This latter assumption is attributed to Mader and White (6) and is referred to by them as a “one-step” approximation. Other assumptions of this model are: (i) the system is isothermal, (ii) dilute solution theory applies, (iii) the flow is laminar in the reactor, and (iv) the Nernst-Einstein relationship applies. The assumption of laminar flow within the reactor could be changed to turbulent flow, if desired. However, this would require that a turbulent velocity profile be used when determining the various average velocities throughout the reactor. It would also require that a turbulent velocity profile be used when determining the velocity-averaged concentrations.

In this model it is assumed that a parallel plate flow cell with a separator can be represented by the five regions shown in Fig. 1. These regions are: the region near the electrode on the anode side; the bulk region on the anode side, the separator region, the bulk region on the cathode side, and the region near the electrode on the cathode side. Additional regions could be added to the model. For example, diffusion regions near the separator/bulk region interfaces could be added. This was not done here because the gradients of the chemical species in these regions would be expected to be small (8). Each region could transfer material to the regions next to it through the processes of diffusion and migration, and each region except the separator has material carried into and out of it by the process of convection. For the purpose of simplicity, the velocity profile within each region is assumed to be uniform, and the averaged velocity for each region is determined by an integral averaging method. Also for simplicity, the outlet concentration profiles for each region are assumed to be uniform, although these values are also determined by an integral-
averaging method. It is also assumed that the potential and concentration profiles within the regions near each electrode and within the separator are linear. Another assumption is that the bulk region of each flow channel is perfectly mixed, and as a result, the concentration profiles within each bulk region are flat. If the perfectly mixed assumption for the bulk region is correct, then it follows that the electrical resistivity in each bulk region is constant. Hence, it follows that the potential gradient within each bulk region must be constant.

Prior to developing the material balance equations for the reactor, it is necessary to determine the average fluid velocity in each of the regions of the reactor. This is done in a standard way in Appendix A (for each diffusion region) and in Appendix B (for each perfectly mixed region) of Ref. (14). The results are as follows:

\[
\bar{v}_d = \frac{\bar{v}_A}{S_A} \left( 3 - 2 \frac{\delta_A}{S_A} \right) \quad [1]
\]

\[
\bar{v}_e = \frac{\bar{v}_C}{S_C} \left( 3 - 2 \frac{\delta_C}{S_C} \right) \quad [2]
\]

for each diffusion region (i.e., within the \(\delta_A\) and \(\delta_C\) regions of Fig. 1) and

\[
\bar{v}_{ae} = \frac{\bar{v}_A}{S_A - \delta_A} \left( S_A - 3 \frac{\delta_A^2}{S_A} + 2 \frac{\delta_A^3}{S_A^2} \right) \quad [3]
\]

\[
\bar{v}_{ec} = \frac{\bar{v}_C}{S_C - \delta_C} \left( S_C - 3 \frac{\delta_C^2}{S_C} + 2 \frac{\delta_C^3}{S_C^2} \right) \quad [4]
\]

for each perfectly mixed region. It is also necessary to determine the average concentrations in each of the regions. These average concentrations are as follows:

\[
\bar{C}_{ae} = \frac{1}{2} \left( C_{a_1} + C_{b_1} \right) \quad [5]
\]

\[
\bar{C}_{ec} = \frac{1}{2} \left( C_{c_1} + C_{b_2} \right) \quad [6]
\]

for the volume-averaged concentrations near the anode

\[
\bar{C}_{ce} = \frac{1}{2} \left( C_{c_1} + C_{c_2} \right) \quad [7]
\]

for the volume-averaged concentrations near the cathode

\[
\bar{C}_{c_1} = \frac{1}{2} \left( C_{c_1} + C_{c_2} \right) \quad [8]
\]

for the volume-averaged concentrations in the separator

\[
\bar{C}_{c_1} = C_{c_1} + (C_{b_1} - C_{c_2}) \left( \frac{2 - 3 \delta_A}{2 \delta_A} \right) \quad [9]
\]

for the velocity-averaged concentrations near the anode, and

\[
\bar{C}_{c_1} = C_{c_1} + (C_{b_1} - C_{c_2}) \left( \frac{2 - 3 \delta_c}{2 \delta_c} \right) \quad [10]
\]

for the velocity-averaged concentrations near the cathode. Equations [5-9] are derived in Ref. (14).

The material balance equations for the five regions between the electrodes are

\[
\frac{d\bar{C}_{ae}}{dt} = \bar{v}_A \frac{N_{a_1} - N_{b_1}}{S_A} + \bar{v}_{ae} \left( C_{a_1} - \bar{C}_{ae} \right) + R_{ae} \quad [11]
\]

(diffusion region near anode)

\[
\frac{d\bar{C}_{ec}}{dt} = \bar{v}_C \frac{N_{c_1} - N_{c_2}}{S_C} + \bar{v}_{ec} \left( C_{c_1} - \bar{C}_{ec} \right) + R_{ec} \quad [12]
\]

(bulk region near anode)

\[
\frac{d\bar{C}_{c_1}}{dt} = \bar{v}_C \frac{N_{c_1} - N_{c_2}}{S_C} + \bar{v}_{c_1} \left( C_{c_1} - \bar{C}_{c_1} \right) + R_{c_1} \quad [13]
\]

(bulk region near cathode)

\[
\frac{d\bar{C}_{c_1}}{dt} = \bar{v}_C \frac{N_{c_1} - N_{c_2}}{S_C} + \bar{v}_{c_2} \left( C_{c_1} - \bar{C}_{c_1} \right) + R_{c_2} \quad [14]
\]

(diffusion region near cathode)

Equations [10]-[14] contain one or more variables representing molar flux. These quantities are defined as follows

\[
\begin{align*}
N_{ai} & = -D_i \frac{C_{ai} - C_{ba}i}{\delta_a} - ziF \frac{D_i}{RT} C_{ai} \phi_i - \phi_a \\
N_{bi} & = -D_i \frac{C_{bi} - C_{ba}i}{\delta_b} - ziF \frac{D_i}{N_MRT} C_{bi} \phi_i - \phi_b \\
N_{ci} & = -D_i \frac{C_{ci} - C_{ba}i}{\delta_c} - ziF \frac{D_i}{RT} C_{ci} \phi_i - \phi_c \\
N_{ds} & = -D_i \frac{C_{ds} - C_{ba}i}{\delta_d} - ziF \frac{D_i}{RT} C_{ds} \phi_i - \phi_d
\end{align*}
\]

Equations [10] and [14] contain one of two variables representing reaction rate. These reaction rates are Butler-Volmer kinetic expressions and are defined as follows

\[
\begin{align*}
\eta_{ai} & = \beta \frac{g_{ai}^0}{{n_iF}} \left( \frac{C_{ai} - C_{ba}i}{C_{ai} - C_{ba}i} \right)^{\alpha_{ai}F \frac{RT}{RT}} \eta_{ai}^{0.5} \left( \frac{C_{ai} - C_{ba}i}{C_{ai} - C_{ba}i} \right)^{\alpha_{ai}F \frac{RT}{RT}} \eta_{ai}^{0.5} \right) \\
\eta_{bi} & = \beta \frac{g_{bi}^0}{{n_iF}} \left( \frac{C_{bi} - C_{ba}i}{C_{bi} - C_{ba}i} \right)^{\alpha_{bi}F \frac{RT}{RT}} \eta_{bi}^{0.5} \left( \frac{C_{bi} - C_{ba}i}{C_{bi} - C_{ba}i} \right)^{\alpha_{bi}F \frac{RT}{RT}} \eta_{bi}^{0.5} \\
\eta_{ci} & = \beta \frac{g_{ci}^0}{{n_iF}} \left( \frac{C_{ci} - C_{ba}i}{C_{ci} - C_{ba}i} \right)^{\alpha_{ci}F \frac{RT}{RT}} \eta_{ci}^{0.5} \left( \frac{C_{ci} - C_{ba}i}{C_{ci} - C_{ba}i} \right)^{\alpha_{ci}F \frac{RT}{RT}} \eta_{ci}^{0.5} \\
\eta_{ds} & = \beta \frac{g_{ds}^0}{{n_iF}} \left( \frac{C_{ds} - C_{ba}i}{C_{ds} - C_{ba}i} \right)^{\alpha_{ds}F \frac{RT}{RT}} \eta_{ds}^{0.5} \left( \frac{C_{ds} - C_{ba}i}{C_{ds} - C_{ba}i} \right)^{\alpha_{ds}F \frac{RT}{RT}} \eta_{ds}^{0.5}
\end{align*}
\]

The charge balance equations for the reactor are

\[
\begin{align*}
\frac{dC_{ai}}{dt} & = 2 \left[ \frac{N_{ai}}{\delta_a} + \frac{\phi_a}{L} (C_{ai} - \bar{C}_{ai}) + R_{ai} \right] - \frac{dC_{ai}}{dt} \\
\frac{dC_{bi}}{dt} & = 2 \left[ \frac{N_{bi}}{\delta_b} + \frac{\phi_b}{L} (C_{bi} - \bar{C}_{bi}) + R_{bi} \right] - \frac{dC_{bi}}{dt} \\
\frac{dC_{ci}}{dt} & = 2 \left[ \frac{N_{ci}}{\delta_c} + \frac{\phi_c}{L} (C_{ci} - \bar{C}_{ci}) + R_{ci} \right] - \frac{dC_{ci}}{dt} \\
\frac{dC_{ds}}{dt} & = 2 \left[ \frac{N_{ds}}{\delta_d} + \frac{\phi_d}{L} (C_{ds} - \bar{C}_{ds}) + R_{ds} \right] - \frac{dC_{ds}}{dt}
\end{align*}
\]

The current density is defined as (10)

\[
i = \frac{F \eta_{ai}^0}{\eta_{ai}^0} = -F \eta_{ai}^0 \eta_{ai}^{0.5}
\]

The voltage drop through the electrolyte during charge is defined as

\[
\Delta \phi = \phi_a - \phi_c
\]

Model Solution

Eq. [28], [11], [13], and [29] govern the time dependence of each of the n species concentrations in each of the four regions (δ_a, δ_b, δ_c, and δ_d) of the reactor. These n material balance equations plus the six current balance equations (i.e., Eq. [30]-[35]) constitute a system of 4n + 6 independent equations. These equations are in terms of 4n + 6 independent concentrations (C_{ai}, C_{bi}, C_{ci}, and C_{ds}) and six potentials (φ_a, φ_b, φ_c, φ_d, φ_e, and φ_f).
Table I. Butler-Volmer kinetic data (T = 298.15 K)

<table>
<thead>
<tr>
<th>Reaction (j) (A-cm²)</th>
<th>$i_{0j}$</th>
<th>$a_{ij}$</th>
<th>$a_{ij}$</th>
<th>$n_j$</th>
<th>$U_{fj}$</th>
<th>$U_{sfej}$</th>
</tr>
</thead>
<tbody>
<tr>
<td>1</td>
<td>+0.31E-02</td>
<td>0.5</td>
<td>0.5</td>
<td>1</td>
<td>+1.087</td>
<td>+1.783</td>
</tr>
<tr>
<td>2</td>
<td>+0.10E+01</td>
<td>0.5</td>
<td>0.5</td>
<td>1</td>
<td>-0.763</td>
<td>+0.000</td>
</tr>
</tbody>
</table>

Table II. Stoichiometry data

<table>
<thead>
<tr>
<th>Species (i)</th>
<th>Reaction 1</th>
<th>Reaction 2</th>
</tr>
</thead>
<tbody>
<tr>
<td>Na⁺</td>
<td>+1</td>
<td>+0.0</td>
</tr>
<tr>
<td>Br⁻</td>
<td>-1</td>
<td>+0.0</td>
</tr>
<tr>
<td>Br₂</td>
<td>+2</td>
<td>+0.5</td>
</tr>
<tr>
<td>Zn²⁺</td>
<td>+2</td>
<td>+0.0</td>
</tr>
<tr>
<td>Br₃⁻</td>
<td>-1</td>
<td>+0.0</td>
</tr>
</tbody>
</table>

Table III. Physical property data and feed composition

<table>
<thead>
<tr>
<th>Species (i)</th>
<th>$z_i$</th>
<th>$D_i$ (cm²/s)</th>
<th>$C_i$ (mol/cm³)</th>
</tr>
</thead>
<tbody>
<tr>
<td>Na⁺</td>
<td>+1</td>
<td>+0.1334E-04</td>
<td>+0.1000E-02</td>
</tr>
<tr>
<td>Br⁻</td>
<td>-1</td>
<td>+0.2084E-04</td>
<td>+0.2949E-02</td>
</tr>
<tr>
<td>Br₂</td>
<td>+2</td>
<td>+0.1310E-04</td>
<td>+0.1015E-05</td>
</tr>
<tr>
<td>Zn²⁺</td>
<td>+2</td>
<td>+0.7540E-05</td>
<td>+0.1000E-02</td>
</tr>
<tr>
<td>Br₃⁻</td>
<td>-1</td>
<td>+0.1310E-04</td>
<td>+0.5150E-04</td>
</tr>
</tbody>
</table>

The selection of values for the diffusion layer thicknesses is critical to this model. It may be possible to estimate values for these thicknesses by using boundary layer theory. It may also be possible to determine these values by experimental measurement. The authors estimated values for these thicknesses by comparing predictions from the model (using assumed values for the thicknesses) with predictions from previous models.

Once the steady-state conditions within the reactor have been predicted, the average outlet concentrations are readily determined by a material balance at the exit of the reactor. The resulting equations are

$$C_{fai} = \frac{\delta_a \theta_{ai} C_{ai} + (S_A - \delta_a) \theta_{bd} C_{fbi}}{S_A \theta_a}$$

$$C_{fci} = \frac{\delta_c \theta_{ci} C_{ci} + (S_C - \delta_c) \theta_{bd} C_{fbi}}{S_C \theta_c}$$

Once the outlet concentrations are determined, the conversions for each of the components can easily be determined from the following definitions

$$x_{ai} = \frac{C_{fai} - C_{ci}}{C_{fai}}$$

$$x_{ci} = \frac{C_{fci} - C_{ai}}{C_{fci}}$$

### Discussion and Results

Figures 2 and 3 are comparisons of predictions obtained from Mader and White (6) and the model presented here. In both of these figures, the solid line represents the predictions obtained from Mader’s “one-step” model, and the dashed line represent values obtained from the algebraic model. The comparison between the two models is good (within about 2%) throughout the entire range of effective separator thicknesses investigated. Table IV presents a direct comparison of predictions made using the algebraic model, Mader’s “one-step” model, and Mader’s continuous
model. Surprisingly, the algebraic model predicts values for $i$ and $\delta$ that are closer to those predicted by the continuous model than does the "one-step" model. However, the value for the total efficiency (see Eq. [41] of Ref. (6)) predicted by the algebraic model differs significantly from the values predicted by both of Mader's models. This discrepancy is most likely due to the tremendous difference in the values of the exchange current densities for the zinc and bromine reactions at the zinc electrode (see Table I).

Figure 4 shows how the predicted current density for the cell depends on various assumed values of the diffusion layer thickness. It is apparent from Fig. 4 that the calculated value of the current density depends on the assumed value of the diffusion layer thickness; however, this dependence is weak (about 8% maximum change). This dependence could be used to the experimental data to fit model predictions by adjusting the value of $\delta$.

Figure 5 shows the effect of $E_{cell}$ upon the current efficiency at the cathode. The current efficiency at the cathode is strongly influenced by the operating potential, and it should approach unity for very high values of $E_{cell}$. These two findings are attributed to the tremendous difference between the values of the exchange current densities for the two reactions at the cathode. Figure 6 shows the effect of $E_{cell}$ upon the total energy efficiency of the cell. It is apparent from Fig. 6 that there is an optimum value for the cell potential at which the battery should be operated so that the maximum amount of energy can be stored. This maximum occurs because the total efficiency is the product of the current efficiency and the voltaic efficiency. Although the current efficiency increases monotonically with respect to $E_{cell}$, the voltaic efficiency decreases monotonically with respect to $E_{cell}$. For this system there is a point beyond which any increase in the current efficiency is more than offset by the associated decrease in the voltaic efficiency. For the feed composition studied, the maximum energy efficiency occurs at about 2.10V. Mader (13) presents figures similar to Fig. 5 and 6. Unfortunately, he does not present data beyond 2.03V. As a result, he did not report that the total efficiency has a maximum value.

Figures 7 and 8 show the effect of varying $\beta$ (i.e., the area ratio) for the porous bromine electrode. It seems reasona-
It is likely that the optimum value of $E_{\text{cell}}$ at which the maximum total efficiency occurs is influenced by the feed composition. If this is true, then it follows that for a system where the feed composition changes with time (such as the zinc/bromine battery), the optimum value of $E_{\text{cell}}$ would change with time. The occurrence of an optimum $E_{\text{cell}}$ is a very practical result that may be applicable to all parallel plate cells with multiple reactions at one or more of the electrodes.

**Conclusions**

This paper demonstrates that a parallel plate flow cell can be approximated by a combination of perfectly mixed regions and diffusion regions. For a reactor that incorporates a porous separator and has a low conversion (such as the zinc/bromine battery), the reactor can be approximated at steady state by $3n + 4$ variables. Also, this paper presents a set of ordinary differential equations that describe the performance of a parallel flow cell at unsteady state. These equations should be solvable by a Runge-Kutta technique, and such a solution may be the subject of a future paper.

This paper also shows that an optimum value of $E_{\text{cell}}$ may exist for a zinc/bromine battery. For the feed composition studied, the optimum cell potential is about 2.10V. Such an optimum may also exist for other systems with multiple electrode reactions. For the zinc/bromine system, the optimum cell potential may be influenced by the feed composition. If this is true, then the value of $E_{\text{cell}}$ would have to be constantly varied for the battery to be operated in the most efficient manner.

**Acknowledgment**

The authors wish to acknowledge the support of this work by the National Science Foundation through Grant No. CBT-8620142.
\[ \frac{r_i - N_{n_i}}{\delta_A} + \frac{\Phi}{D_i} \left( C_{ba} - C_{ba}^* \right) + R_{ technically necessary} \]

\[ = \frac{C_{ba}(t + \Delta t) - C_{ba}(t)}{\Delta t} \quad [A-1.2] \]

Taking the limit of Eq. [A-1.2] as \( \Delta t \) approaches zero yields

\[ \frac{r_i - N_{n_i}}{\delta_A} + \frac{\Phi}{D_i} \left( C_{ba} - C_{ba}^* \right) + R_{ technically necessary} = \frac{dC_{ba}}{dt} \]

**APPENDIX B**

Figure A-2 is a schematic of the mole fluxes that enter and leave the bulk region on the anode side.

\[ N_{n_i} = -D_i \frac{C_{ba} - C_{ba}^*}{\delta_A} \frac{C_{ba} - C_{ba}^*}{\delta_A} \frac{C_{ba} - C_{ba}^*}{\delta_A} \frac{C_{ba} - C_{ba}^*}{\delta_A} \]

\[ N_{n_i} = -D_i \frac{C_{ba} - C_{ba}^*}{\delta_A} \frac{C_{ba} - C_{ba}^*}{\delta_A} \frac{C_{ba} - C_{ba}^*}{\delta_A} \frac{C_{ba} - C_{ba}^*}{\delta_A} \]

A verbal statement of the material balance is

\[ \text{(moles in)} - \text{(moles out)} + \text{(moles generated)} = \text{(moles accumulated)} \]

where

\[ \text{(moles in)} = [N_{n_i} L W + C_{ba} \delta_A] \frac{!}{!} \]

\[ \text{(moles out)} = [N_{n_i} L W + C_{ba} \delta_A] \frac{!}{!} \]

\[ \text{(moles generated)} = R_{ technically necessary} \]

\[ \text{(moles accumulated)} = [C_{ba}(t + \Delta t) - C_{ba}(t)] L \delta_A \]

Let \( S = S_A - \delta_A \). Substitution of these definitions into the verbal material balance yields

\[ N_{n_i} L W + C_{ba} \delta_A L S - (N_{n_i} L W + C_{ba} \delta_A L S) \Delta t \]

\[ = (C_{ba}(t + \Delta t) - C_{ba}(t)) L S \]

\[ [A-2.1] \]

Rearrangement of Eq. [A-2.1] yields the following

\[ \frac{N_{n_i} - N_{n_i}}{S_A - \delta_A} + \frac{\Phi}{D_i} \left( C_{ba} - C_{ba}^* \right) + R_{ technically necessary} \]

\[ \frac{\Phi}{D_i} \left( C_{ba} - C_{ba}^* \right) + R_{ technically necessary} = \frac{C_{ba}(t + \Delta t) - C_{ba}(t)}{\Delta t} \quad [A-2.2] \]
Fig. A-3. Charge balance schematic of the diffusion region near the anode.

Taking the limit of Eq. [A-2.2] as $\Delta t$ approaches zero yields the following

\[
\frac{N_{b_i} - N_{b_{in}}}{S_A - \delta_A} - \frac{D_i}{L} (C_{b_{in}} - C_{b_i}) + R_{b_{in}} = \frac{dC_{b_{in}}}{dt}
\]

APPENDIX C

Figure A-3 is a schematic of the charge fluxes that enter and leave the diffusion region near the anode

\[
N_{b_i} = -D_i \frac{C_{b_{in}} - C_{b_i}}{\delta_A} - \frac{z_F D_i}{RT} \frac{C_{b_{in}} \phi_i - \phi_1}{\delta_A}
\]

A verbal statement of the charge balance is

\[
\text{(charge in)} - \text{(charge out)} = 0
\]

where

\[
\text{(charge in)} = FLW \Sigma z_i r_i
\]

\[
\text{(charge out)} = FLW \Sigma z_i N_{b_i}
\]

Substitution of these definitions into the verbal charge balance yields

\[
FLW \Sigma z_i r_i - FLW \Sigma z_i N_{b_i} = 0 \quad \text{[A-3.1]}
\]

Simplification and substitution of the flux expression into Eq. [A-3.1] yield

\[
\Sigma z_i (r_i - \Sigma z_i) \left( -D_i \frac{C_{b_{in}} - C_{b_i}}{\delta_A} - \frac{z_F D_i}{RT} \frac{C_{b_{in}} \phi_i - \phi_1}{\delta_A} \right) = 0
\]

A rearrangement of Eq. [A-3.2] yields the following

\[
\Sigma z_i r_i + \frac{1}{\delta_A} \left[ \Sigma z_i D_i (C_{b_{in}} - C_{b_i}) + \frac{F(\phi_i - \phi_1)}{RT} \Sigma z_i^2 D_i C_{b_{in}} \right] = 0
\]

APPENDIX D

Figure A-4 is a schematic of the charge fluxes that enter and leave the bulk region on the anode side

\[
N_{b_{in}} = -z_F \frac{D_i}{RT} \frac{C_{b_{in}} \phi_i - \phi_1}{S_A - \delta_A}
\]

A verbal statement of the charge balance is

\[
\text{(charge in)} - \text{(charge out)} = 0
\]

where

\[
\text{(charge in)} = FLW \Sigma z_i r_i
\]

\[
\text{(charge out)} = FLW \Sigma z_i N_{b_{in}}
\]

Substitution of these definitions into the verbal charge balance yields

\[
FLW \Sigma z_i r_i - FLW \Sigma z_i N_{b_{in}} = 0 \quad \text{[A-4.1]}
\]

Substituting the flux expression into Eq. [A-4.1] yields

\[
\Sigma z_i r_i - \Sigma z_i \left( -z_F \frac{D_i}{RT} \frac{C_{b_{in}} \phi_i - \phi_1}{S_A - \delta_A} \right) = 0
\]

A rearrangement of Eq. [A-4.2] yields the following

\[
\Sigma z_i r_i + \frac{F(\phi_2 - \phi_1)}{RT(S_A - \delta_A)} \Sigma z_i^2 D_i C_{b_{in}} = 0
\]

REFERENCES