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## An Infrared Imaging Method for High-Throughput Combinatorial Investigation of Hydrogenation-Dehydrogenation and New Phase Formation of Thin Films

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# An infrared imaging method for high-throughput combinatorial investigation of hydrogenation-dehydrogenation and new phase formation of thin films

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We have developed an infrared imaging setup enabling *in situ* infrared images to be acquired, and expanded on capabilities of an infrared imaging as a high-throughput screening technique, determination of a critical thickness of a Pd capping layer which significantly blocks infrared emission from below, enhancement of sensitivity to hydrogenation and dehydrogenation by normalizing raw infrared intensity of a Mg thin film to an inert reference, rapid and systematic screening of hydrogenation and dehydrogenation properties of a Mg–Ni composition spread covered by a thickness gradient Pd capping layer, and detection of formation of a Mg<sub>2</sub>Si phase in a Mg thin film on a thermally oxidized Si substrate during annealing. © 2009 American Institute of Physics. [DOI: 10.1063/1.3184024]

## I. INTRODUCTION

Due to the pressing social need for alternative sources of energy, especially for transportation applications, the search for new lightweight hydrogen storage materials has intensified.<sup>1,2</sup> The ideal material would take on a large amount of hydrogen per weight and volume at a moderate temperature in the time it would take to fill up a conventional car's gas tank. To date, the majority of this research has been done using traditional one-by-one synthesis and measurement techniques.<sup>3–16</sup>

Although a combinatorial approach, where a large number of samples are synthesized in a single experiment, processed, and then rapidly characterized in parallel,<sup>17–23</sup> has the potential to significantly expedite the discovery of new hydrogen storage materials, only a few groups have adopted it so far.<sup>24–28</sup> In large part, the hesitance to this new methodology is based on the lack of an effective screening technique, frequently the bottleneck for new combinatorial studies. A few screening techniques for hydrogen storage materials have been proposed such as deflection of cantilever arrays,<sup>24</sup> nanocalorimetry,<sup>25</sup> and an optical imaging.<sup>26,27</sup> In addition, infrared (IR) imaging put forward by Olk *et al.*<sup>28</sup> is one of the promising screening techniques. This method can be complementary to the optical imaging as some materials are active for IR but not for visible light and vice versa due to their different properties in visible and IR spectral regions.

In thermodynamics, Kirchhoff's law of thermal radiation is a general statement equating emission and absorption in

heated objects, as proposed by Gustav Kirchhoff in 1859.<sup>29</sup> Accordingly, normal spectral emissivity  $\varepsilon(\lambda)$  is related to normal spectral reflectivity  $R(\lambda)$  as

$$\varepsilon(\lambda) = 1 - R(\lambda). \quad (1)$$

Reflectivity and emissivity in the IR are linked to electrical conductivity via the complex dielectric function. Reflectivity (for normal incidence) depends on the property of material through components of its complex refractive index  $N = n + i\kappa$

$$R = \frac{(n - n_0) + \kappa^2}{(n + n_0) + \kappa^2}, \quad (2)$$

where  $n$  is the refractive index,  $\kappa$  is the extinction coefficient, and  $n_0$  is the refractive index of the ambient medium. From this, the relationship between emissivity and conductivity can be established.<sup>30–32</sup>

In the IR spectral region where the wavelength of the light is sufficiently long, the Hagen–Rubens relation (for normal incidence of radiation) provides a direct relationship between emissivity and electrical resistivity  $\rho$  (Refs. 30 and 31)

$$\varepsilon(\omega) \approx 2\sqrt{2\varepsilon_0\omega\rho}, \quad (3)$$

where  $\varepsilon_0$  is the dielectric constant of vacuum and  $\omega$  is the angular frequency of electromagnetic wave.

From these relations, it follows that metals generally have considerably lower emissivities than insulators. The absorption of hydrogen in metals and alloys results in new scattering centers and eventually leads to a modified density of states, which is often accompanied by a rise in the Fermi energy  $E_f$ .<sup>33–36</sup> The shift in  $E_f$  to an energy with a lower density of states leads to a less metallic character and there-

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fore an increase in the electrical resistance or change in conductivity from metallic to insulating. Thus IR emissivity is expected to increase with the increase in hydrogen in a material.

Surface roughness of the material also affects emissivity. In the IR spectral region, reflectivity is expressed as

$$R = R_0 \exp[-(4\pi\sigma)^2/\lambda^2], \quad (4)$$

where  $R_0$  is the specular reflectivity of a perfectly smooth surface and  $\sigma$  is the root mean square roughness of the surface.<sup>37</sup> However, considering microscopically smooth surface after hydrogenation,<sup>38</sup> the effect of the surface roughness is expected to be negligible in our study.

Every material is a natural IR emitter at ordinary temperature range, thus no light source is necessary for imaging measurement. In addition, as IR imaging is sensitive to temperature change because of IR emission's proportionality to the fourth power of temperature,<sup>29</sup> it has a capability to detect the tiny temperature change due to heat of reaction corresponding to hydrogenation and dehydrogenation.

In this paper, the samples were prepared in the form of a thin film, whose unique advantages sometimes work effectively on the investigation of the hydrogen storage materials. For example, superior control on the layer thickness and interface was applied to make multilayered thin films, which improved hydrogenation and dehydrogenation properties.<sup>39,40</sup> Thin films also enabled sophisticated study of the hydrogen reaction on the Mg surface layer.<sup>41,42</sup> In spite of these advantages, one always has to keep it in his mind that large stresses generated in thin films due to clamping to the substrate changes hydrogenation and dehydrogenation properties, and simple comparison to bulk samples becomes difficult.<sup>43,44</sup> However, thin films still insist on their benefits in the combinatorial approach for the materials' discovery since properties dependence on systematically varied factors, such as compositions, can be compared with each other under homogeneously induced stress.<sup>26</sup>

## II. EXPERIMENTAL SETUP

Shown in Fig. 1(a) is a top view diagram of the IR imaging setup. Hydrogenation and dehydrogenation of the samples take place inside the vacuum reaction chamber (red) with six 11.43 cm diameter conflat flanges (two are not used and not shown in the figure).

Working counterclockwise from the left, the first flange is connected to normal purity Ar gas (99.9%) and ultrahigh purity H<sub>2</sub> gas (99.999%) for purging and hydrogenation. The gas pressure is measured by a pressure gauge. For safety, the extra pressure is released from a back-pressure regulator.

The second is used to mount a heating stage. The thin film samples are mounted on the heating stage electronically connected to a temperature controller. Figure 1(b) shows a top view diagram of the heating stage. In the figure, as an example, a Mg thin film sample deposited on a substrate and a  $k$ -type thermocouple are shown clamped on the heating stage by stainless steel clamps, which are screwed into the edge of the stage. The heating stage has two concentric frames of 2.79 and 3.43 cm diameters. A gap between the inner and outer frames increases the heating efficiency. The

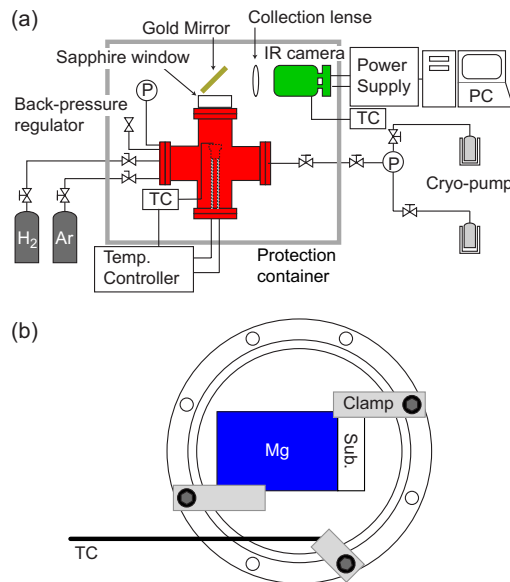


FIG. 1. (Color online) (a) A top view diagram of the IR imaging setup. The reaction chamber with the heating stage and the IR camera are in red and green, respectively. TC, P, and PC denote the  $k$ -type thermocouple, the pressure gauge, and the personal computer, respectively. (b) A top view diagram of the heating stage.

heating stage can withstand up to 1473 K in vacuum and its temperature is controlled within 1 K accuracy.

The third connects to a vacuum system. Liquid nitrogen molecular sieve dual-stage cryopumps are used to evacuate the chamber.

The fourth offers optical imaging of the sample through a 4.45 cm diameter by a 0.63 cm thick sapphire window facing parallel to the heating stage. Rubber o-rings (not shown) on both sides of the window ensure sealing, and protect the window from point contacts to metal. During imaging measurements, IR emission between wavelength of 0.1 and 5.5  $\mu\text{m}$  from the thin film sample and the substrate is transmitted through the sapphire window, reflected by a gold mirror, and imaged onto the focal plane of an IR detector via a transparent CaF<sub>2</sub> collection lens. A micrometer (not shown) attached to the base of the camera permits fine-tuning of the distance between the collection lens and the IR detector to optimize focusing of the sample image. Double-sided polished Si wafers (not shown) placed between the gold mirror and the collection lens serve as 50% attenuators to prevent saturation of the detector.

The liquid-nitrogen-cooled IR detector was used for low noise operation. The temperature of the detector is monitored by a  $k$ -type thermocouple directly inserted into the camera Dewar. The IR detector consists of an array of  $256 \times 256$  InSb diode pixels and dual output video is internally triggered at 25 Hz (or slower) to permit "snap-shot" frame imaging with 10 ms or longer integration times. Custom-made software periodically collects the wavelength-integrated IR emission, which is amplified by multistage analog amplifiers attached to the camera. Analog data is digitally converted by dual 16-bit A/D converters and the image is processed by a digital frame-grabber (Mutek).<sup>45</sup> The program then collects each IR image consisting of a  $256 \times 256$  IR intensity data matrix. The image collection rate per time is controlled by an

external trigger source that can vary from 40 ms to several minutes. Each image is then saved as a binary image file for later processing and image analysis.

The entire system is set on a  $76.20 \times 67.31 \times 6.35$  cm<sup>3</sup> breadboard. For safety from potential high pressure explosion, the system is enclosed in a  $76.20 \times 85.09 \times 101.60$  cm<sup>3</sup> protection container using high-impact, safety-rated plastic windows with two doors, which allow access to the chamber during sample loading.

### III. DEMONSTRATION

All the thin film samples discussed in this paper were deposited in a combinatorial electron beam deposition chamber at room temperature (RT). Details of the deposition schemes for composition spread and thickness gradient thin film samples can be found elsewhere.<sup>46,47</sup> Prior to the IR imaging, all samples were transferred in ambient atmosphere from the deposition chamber to the reaction chamber. After loading the samples, the reaction chamber was Ar purged twice and pumped down to  $1 \times 10^{-6}$  bar by the cryopumps.

Throughout the IR imaging experiments, IR images were collected every 30 s. In differential imaging mode, which was exclusively used in this study, an IR image collected at RT was used as the background; The raw IR intensity is defined as [(IR intensity at measurement temperature) – (IR intensity at RT)].

#### A. Effect of Pd capping layer on IR imaging

Pd is commonly used in hydrogen storage thin film experiments as a capping layer to prevent oxidation and to facilitate splitting of hydrogen molecules at the Pd surface.<sup>48–50</sup> However, metallic Pd with small electrical resistivity  $\rho$  and large reflectivity  $R$  [see Eqs. (1) and (3)] is expected to block IR emission from below. Thus if one is interested in studying the hydrogenation and dehydrogenation of the materials under the Pd capping layer, one has to know how the thickness of the Pd capping layer affects the IR imaging.

To study the effect of the Pd capping layer on the IR imaging, we recorded the IR images of a Pd (0–50 nm)/C (carbon) (300 nm) bilayer thickness gradient thin film sample on an Al<sub>2</sub>O<sub>3</sub> (0001) substrate [Fig. 2(a)]. During measurement, the sample temperature ( $T_s$ ) was increased from 323 to 423 K at 5 K/min and then kept at 423 K for 70 min. Carbon was chosen due to its high IR emissivity.

Figure 2(b) shows a representative top view IR image of the sample recorded at 423 K. The position of the clamp is outlined by the dashed white line. In order to analyze the effect of the Pd capping layer on the IR imaging, the raw IR intensities extracted from the IR image are plotted as a function of the Pd thickness ( $t_{Pd}$ ) in Fig. 2(c). Up to  $t_{Pd} \approx 8$  nm the raw IR intensity increases, and then for  $t_{Pd} > 8$  nm it gradually decreases until around 30 nm. The constant intensity for more than  $t_{Pd} = 30$  nm indicates that the IR emission from the carbon layer is completely blocked by the metallic Pd layer if its thickness exceeds  $\approx 30$  nm. This result provided guidance for designing thin film samples in the following experiments.

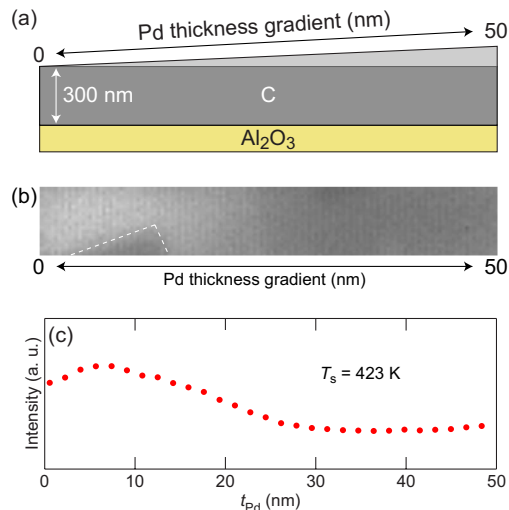


FIG. 2. (Color online) (a) A cross-sectional schematic and (b) a top view IR image of a Pd/C bilayer thickness gradient sample at 423 K deposited on an Al<sub>2</sub>O<sub>3</sub> (0001) substrate. (c) 29 raw IR intensities as a function of the Pd thickness ( $t_{Pd}$ ).

#### B. Normalization of raw IR intensity and detection of hydrogenation and dehydrogenation

Here first importance of normalization of the raw IR intensity is demonstrated in hydrogenation of a 100 nm thick Mg thin film sample deposited on an Al<sub>2</sub>O<sub>3</sub> (0001) substrate [Fig. 3(a)], and second, a capability of the IR imaging to detect hydrogenation and dehydrogenation of the thin film samples is verified.

In the hydrogenation process,  $T_s$  was maintained at 373 K and hydrogen gas pressure ( $P_H$ ) was rapidly increased to 1 bar from vacuum. Obtained IR intensity curves, composed of a large number of raw IR intensities plotted as a function of  $t$ , of the Al<sub>2</sub>O<sub>3</sub> substrate (light gray), and the Mg thin film (dark gray), are shown together with a normalized IR intensity curve (green) [(raw IR intensity curve of Mg)/(raw IR intensity curve of Al<sub>2</sub>O<sub>3</sub>)] in Fig. 3(b). Mg shows significant increase in intensity including a small spike immediately after exposing the sample to hydrogen gas. As Al<sub>2</sub>O<sub>3</sub>, which is inert to hydrogen gas, also shows such a spike, it can be noise. The noise complicates and can mislead data interpretation. On the other hand, normalization smoothes out such an unfavorable feature. Thus, data interpretation becomes easier and more accurate, and the sensitivity to subtle changes in IR emissivity due to hydrogenation and dehydrogenation is remarkably enhanced. Hereafter, the normalization IR intensity curve is used instead of the raw IR intensity curve.

After hydrogenation,  $T_s$  was first reduced to RT and the hydrogen gas was pumped out of the chamber. Then, in the dehydrogenation process, the substrate was heated back up to 423 K at 5 K/min and kept constant. An intensity of an IR intensity curve [Fig. 3(c)] decreases with  $T_s > 400$  K.

Since only hydrogen gas is in and out of the chamber in the hydrogenation and dehydrogenation process, the increase and decrease in IR intensities are presumably attributed to hydrogenation and dehydrogenation of the thin film sample, respectively.

In order to empirically evidence this presumption,  $\theta$ -2 $\theta$



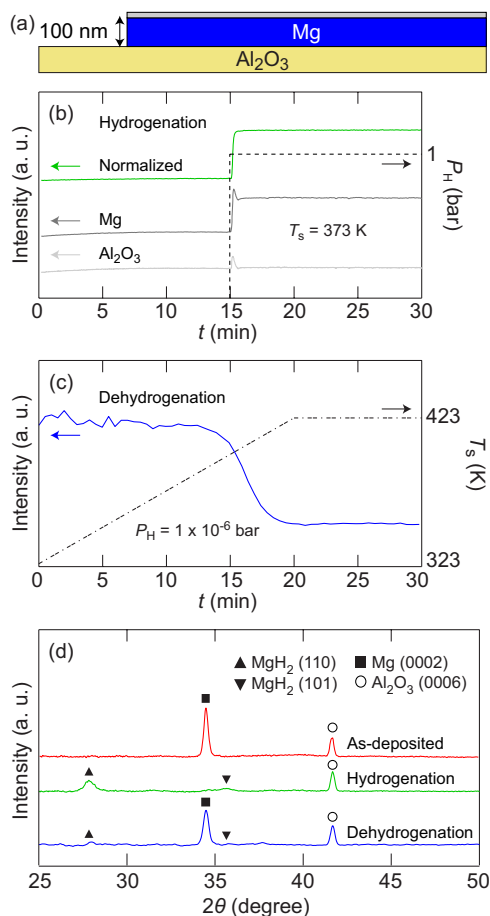


FIG. 3. (Color online) (a) A cross-sectional schematic of a 100 nm thick Mg thin film sample with a 5 nm thick Pd capping layer deposited on an  $\text{Al}_2\text{O}_3$  (0001) substrate. (b) Raw IR intensity curves of the  $\text{Al}_2\text{O}_3$  substrate (light gray) and the Mg thin film (dark gray), and a normalized IR intensity curve (green) obtained in the hydrogenation process. Hydrogen gas pressure ( $P_H$ ) (dashed line) is added for reference (right axis). (c) A normalized IR intensity curve obtained in the dehydrogenation process (left axis) and a sample temperature for reference (chain line) (right axis). (d)  $\theta$ - $2\theta$  XRD (Cu  $K\alpha$  radiation) patterns of a Mg thin film sample before hydrogenation (as-deposited) (red), after hydrogenation (green), and after dehydrogenation (blue).

x-ray diffraction (XRD) was implemented. In Fig. 3(d), other than the  $\text{Al}_2\text{O}_3$  (0006) substrate peak at  $41.7^\circ$ , only the Mg (0002) peak at  $34.4^\circ$  was observed for the as-deposited sample (red). After hydrogenation process (green), the Mg peak almost disappeared. Instead, the  $\text{MgH}_2$  (110) and (101) peaks were observed at  $27.9^\circ$  and  $35.7^\circ$ , respectively. After the dehydrogenation process (blue), the  $\text{MgH}_2$  peaks almost disappeared and the Mg peak reappeared. The reactions suggested by the  $\theta$ - $2\theta$  XRD patterns are only the  $\text{MgH}_2$  formation and decomposition in the hydrogenation and dehydrogenation processes, respectively. This supports our presumption, e.g., the increase and decrease in IR intensities should be directly attributed to hydrogenation and dehydrogenation, respectively.

### C. Systematic investigation of hydrogenation and dehydrogenation in a Mg–Ni composition spread

To illustrate the capability of the IR imaging technique to screen hydrogenation and dehydrogenation properties of a

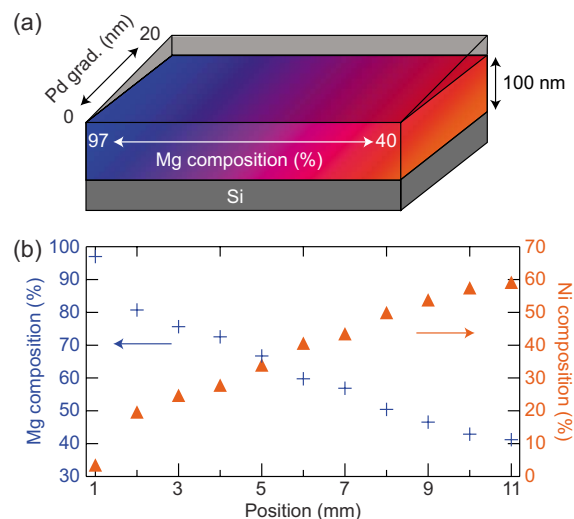


FIG. 4. (Color online) (a) Schematic of a 100 nm thick Mg–Ni composition spread sample deposited on a thermally oxidized Si (100) substrate with a thickness gradient Pd capping layer. (b) Mg–Ni composition variation measured by EDS plotted as a function of position of the sample along the composition spread. Cross and triangle are for the Mg composition (left axis) and for the Ni composition (right axis), respectively.

sample with locally varied parameters, a Mg–Ni composition spread was monitored while cycling hydrogen in and out of the chamber. Figure 4(a) shows the geometry of the Mg–Ni composition spread. In a horizontal direction, a 100 nm thick continuous Mg–Ni composition spread was grown. The composition variation was measured by energy dispersive spectroscopy (EDS) as shown in Fig. 4(b). In the orthogonal direction, a Pd thickness gradient from 0 to 20 nm was grown as a capping layer. This unique sample geometry allows for the simultaneous study of the effect of Mg–Ni composition and thickness of the Pd capping layer on hydrogenation and dehydrogenation properties.

Figures 5(a)–5(e) show representative top view IR images of the Mg–Ni composition spread. Under 1 bar hydrogen gas pressure at  $t=1.5$  h [Fig. 5(b)], the Mg rich side

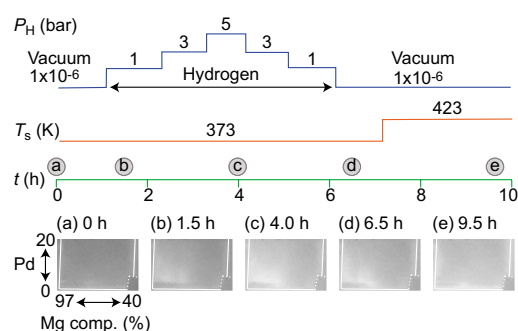


FIG. 5. (Color online) The measurement schedule [hydrogen gas pressure ( $P_H$ ), sample temperature ( $T_s$ ), and time ( $t$ )] and five representative top view IR images of the Mg–Ni composition spread sample at (a) 0 h, (b) 1.5 h, (c) 4.0 h, (d) 6.5 h, and (e) 9.5 h. Initially the sample was stabilized at 373 K in vacuum for 1 h, after which  $P_H$  was increased to 1, 3, and 5 bar every hour. Subsequently  $P_H$  was decreased using the same schedule down to vacuum. Finally,  $T_s$  was increased to 423 K. The image acquisition time is shown on the top of each image and is marked by circle on the time line. The position of the thin film sample and the clamp are outlined by the solid and dashed white lines, respectively. Mg composition and Pd thickness gradient are shown in (a).

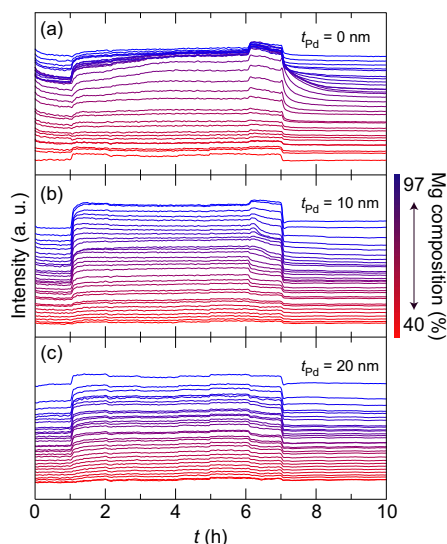


FIG. 6. (Color online) 30 IR intensity curves of measurement spots along the Mg–Ni composition spread for three different thickness of the Pd capping layer ( $t_{\text{Pd}}$ ).  $t_{\text{Pd}}$  (a) 0 nm, (b) 10 nm, and (c) 20 nm, respectively. A Si substrate was used as an inert reference for normalization. The composition of each curve is represented by the curve color as illustrated by the color bar on the right-hand side. The curves are offset vertically for clarity, with the more Mg results placed at the upper side of the figure.

(left) of the sample increases brightness. The increase in brightness is more obvious in the region with the thinner Pd capping layer (bottom), which is consistent with our previous study of the effect of thickness of the Pd capping layer on the IR imaging. Apparently, transformation of Pd to Pd hydride during the experiment does not affect the IR-blocking behavior of the Pd capping layer, probably because of the metallic character of Pd hydride.<sup>51</sup> As  $P_{\text{H}}$  increases, both the magnitude and the spatial range of the brightness increase due to further hydrogenation of the sample [Fig. 5(c)]. In contrast, as  $P_{\text{H}}$  decreases, both the magnitude and the spatial range of the brightness decrease due to partial dehydrogenation of the sample [Fig. 5(d)]. Finally at  $T_{\text{s}}=423$  K, brightness disappeared maybe due to complete dehydrogenation [Fig. 5(e)]. As illustrated here, looking at the IR images enables us to grasp the hydrogenation and dehydrogenation feature of a thin film sample with locally varied parameters immediately.

However, there are some features we cannot understand by simply looking at the IR images. Thus the IR images were analyzed and IR intensity curves are plotted in Figs. 6(a)–6(c). Regardless of Mg–Ni composition and  $t_{\text{Pd}}$ , there are rapid hydrogenation at  $t=1$  h and rapid dehydrogenation at  $t=7$  h indicated by rapid intensity increase and decrease, respectively. In addition, slight dehydrogenation is indicated by a small intensity decrease between  $t=6$  and 7 h, where  $P_{\text{H}}=1 \times 10^{-6}$  bar and  $T_{\text{s}}=373$  K (see schedule in Fig. 5). The remarkable difference among different  $t_{\text{Pd}}$ s is that, only for  $t_{\text{Pd}}=0$  nm, rapid intensity changes are followed by slow intensity changes indicating slow hydrogenation or dehydrogenation for composition range of  $51 \leq \text{Mg composition } (\%) \leq 77$ .

Borgschulte *et al.*<sup>52</sup> previously reported that the kinetics of hydrogenation of the yttrium thin film increases for thicker Pd capping layers. Their result may help explain the

slow increase and decrease in IR intensity observed here for  $t_{\text{Pd}}=0$  nm. Another possible reason is potential oxidation of the thin film surface. Although Mg is very sensitive to oxidation,<sup>53</sup> our sample was transferred in ambient atmosphere from the deposition chamber to the reaction chamber. Thus, the surface with a very thin (and can be discontinuous) Pd capping layer is likely to be partially or fully oxidized. The oxidation would be more severe and suppresses hydrogenation and dehydrogenation in the region with the higher Mg composition. This may explain smaller increase and decrease in IR intensity for  $t_{\text{Pd}}=0$  nm [Fig. 6(a)] than for  $t_{\text{Pd}}=10$  nm [Fig. 6(b)] and 20 nm [Fig. 6(c)] with a composition range between  $80 \leq \text{Mg composition } (\%) \leq 97$ .

In Figs. 6(b) and 6(c), the magnitude of the IR intensity change increases almost linearly with Mg composition in accordance with the increase in hydrogen storage capacity,<sup>54</sup> which suggests the semiquantitative capability of the IR imaging to hydrogen storage capacity. One limitation of the IR imaging is its limited signal-to-noise ratio even after normalization discussed in this report. For example, discontinuous intensity changes occur in Figs. 6(a)–6(c) at  $t \approx 6$  h when  $P_{\text{H}}$  was pumped down to  $1 \times 10^{-6}$  bar. The reason for this discontinuity can be insufficient heating or cooling of the chamber by introduction and mixing of hydrogen gas. Unfortunately, it is difficult to separate the noise from the meaningful intensity change due to hydrogenation and dehydrogenation. In order to discern meaningful information from noise, however, one can plot all the IR intensity curves obtained under the same experimental conditions together using the combinatorial approach. Then systematic intensity changes with the systematically varied physical parameters, e.g., composition, can be delineated.

## D. Detection of formation of a $\text{Mg}_2\text{Si}$ phase

Here we demonstrate capability of the IR imaging to detect formation of new phases by monitoring a Mg thickness gradient thin film sample [Fig. 7(a)] annealed at 573 K in vacuum. Representative top view IR images [Figs. 7(b)–7(d)] show that during the course of annealing, the bright region emerges from the thinner Mg side (left) and propagates toward the thicker Mg side (right). To study details of the behavior of the IR intensity change, IR intensity curves with continuously varied Mg thickness are plotted in Fig. 7(e). A pronounced increase in IR intensity is observed for each curve but it gradually delays with increasing thickness of the Mg thin film. As annealing was implemented in vacuum, the origin of the increase in intensity cannot be attributed to reaction of the Mg thin film with residual gases but rather to some structural changes affecting the intrinsic IR emissivity.

The most plausible explanation for this observation is the interdiffusion of different components of the heterogeneous sample and potential formation of new phases; either Pd diffusion from the capping layer or Si diffusion from the substrate into the Mg thin film, or both. If the change originates from Mg–Pd phase formation due to the Pd diffusion, the diffusion is expected to occur uniformly from the surface and thus the increase in intensity would have been observed simultaneously for each thickness of Mg. This disagrees with

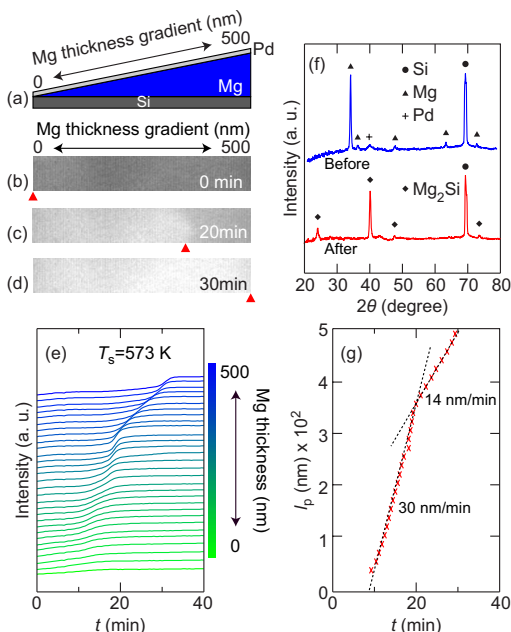


FIG. 7. (Color online) (a) A cross-sectional schematic of a Mg thickness gradient thin film sample with a 5 nm thick Pd capping layer deposited on a thermally oxidized Si (100) substrate. Representative top view IR images of the Mg thickness gradient thin film sample annealed for (b) 0 min, (c) 20 min, and (d) 30 min at 573 K, respectively. A red triangle denotes position of the boundary line between the dark region and bright region. (e) 30 IR intensity curves of measurement spots along the Mg thickness gradient. A Si substrate was used as an inert reference for normalization. The Mg thickness of the IR intensity curves are represented by the curve color as illustrated by the color bar on right-hand side, and offset vertically for clarity with the thicker Mg results placed at the upper side of the figure. (f) Representative  $\theta$ -2 $\theta$  XRD patterns of 500 nm thick Mg before (blue) and after (red) annealing. The diffraction peaks of the Pd and the Mg<sub>2</sub>Si at around 40.3° overlapped and are hard to distinguish. (g) Propagation length ( $l_p$ ) of the Mg<sub>2</sub>Ni phase formation front as a function of time ( $t$ ).

our observation. On the other hand, if the origin is Mg–Si phase formation due to Si diffusion into the Mg thin film, the diffusion would have occurred from the bottom of the Mg thin film, and the increase in intensity would have not been observed until a newly formed Mg–Si phase reached the thin film surface. In this case, the increase in intensity associated with formation of the Mg–Ni phase should be delayed for thicker Mg. Thus the observation supports Si diffusion.

In order to verify formation of the Mg–Si phase, the Mg thickness gradient thin film sample was investigated using XRD before and after annealing. The XRD patterns [Fig. 7(f)] show diffraction peaks originating from Mg, Pd, and Si before annealing (blue) and from Mg<sub>2</sub>Si and Si after annealing (red). This is supportive evidence that significant diffusion of Si has occurred during annealing yielding the formation of a Mg<sub>2</sub>Si phase. As the Mg<sub>2</sub>Si phase is a  $n$ -type semiconductor, the change in the electric characteristic of the thin film from metallic to semiconducting may results in increase in IR emissivity.<sup>55,56</sup>

Provided that the formation front of the Mg<sub>2</sub>Si phase propagates from a substrate toward a thin film surface, the increase in IR intensity begins at the moment the front reached the thin film surface after propagating the distance equals to the thickness of the Mg thin film. The speed of propagation of the Mg<sub>2</sub>Si phase formation front is deter-

mined from the plot of propagation length  $l_p$  (=thickness of the Mg thin film) as a function of  $t$  at which increase in IR intensity begins [Fig. 7(g)]. Apparently there are two slopes in Fig. 7(g); the formation front of the Mg<sub>2</sub>Si phase initially propagates fast, and then slows down. From the slopes of the straight lines on which data points fit, the speed of propagation of the Mg<sub>2</sub>Si phase formation front was determined as 30 and 14 nm/min for fast and slow propagation, respectively. The slowdown may be explained by the formation of the Mg<sub>2</sub>Si continuous layer on the Si substrate, in which Si diffusion might be hindered. Details of the propagation mechanism are under investigation.

#### IV. CONCLUSION

We have developed an IR imaging setup enabling *in situ* IR images to be acquired and expanded on capabilities of an IR imaging as a high-throughput screening technique. The critical thickness of the Pd capping layer, which significantly blocks IR emission from below was experimentally determined to be  $\approx 30$  nm. Normalization of raw IR image data to an inert reference such as a substrate enhances sensitivity to hydrogenation and dehydrogenation. The IR imaging of a Mg thin film in hydrogenation and dehydrogenation process confirmed that an increase and decrease in IR intensity are attributed to hydrogenation and dehydrogenation, respectively. The IR imaging of a Mg–Ni composition spread sample with a thickness gradient Pd capping layer revealed that the amount of absorbed hydrogen changes systematically with composition, and the kinetics changes with the thickness of the Pd capping layer. Finally we have demonstrated an interesting capability of the IR imaging, i.e., detection of phase formation and determination of its propagation speed, by detecting formation of the Mg<sub>2</sub>Si phase in the Mg thin film on the Si substrate during annealing.

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