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## ON THE ATOMIC DECOMPOSITION OF $H^1$ AND INTERPOLATION

ROBERT SHARPLEY<sup>1</sup>

In [1] Coifman used the Fefferman-Stein theory of  $H^p$  spaces [4] to decompose the functions of these spaces into basic building blocks (atoms), further clarifying their real variable nature. Coifman and Weiss have provided a comprehensive treatment of these ideas and many applications to harmonic analysis in [2]. In this note, we use the nontangential maximal function  $Nf$  to give an elementary proof of the decomposition of  $H^1$  functions on the line and then characterize the Peetre  $K$ -functional for  $H^1$  and  $L^\infty$  in terms of  $Nf$ .

Let  $u$  be the harmonic extension [5] of  $f$  to the upper half plane  $\mathbf{R}_+^2$ . For  $x \in \mathbf{R}$ , denote by  $\Gamma_x = \{(z, y) \in \mathbf{R}_+^2 : |x - z| \leq y\}$  the cone with vertex at  $x$ . The nontangential maximal function of  $f$  is defined by  $Nf(x) = \sup\{|u(z, y)| : (z, y) \in \Gamma_x\}$ . We define the (real)  $H^1$  norm of  $f$  to be the standard  $H^1$  norm of  $u + iv$ , where  $v$  is the harmonic conjugate of  $u$  which satisfies  $v(0) = 0$ . A classical result of Hardy and Littlewood asserts that  $\|Nf\|_{L^1} \leq c\|f\|_{H^1}$ . For an interval  $I$  an  $H^1$ -atom is any function  $a_I$  such that  $\int a_I = 0$  and  $|a_I| \leq |I|^{-1}\chi_I$  a.e.

**PROPOSITION.**<sup>2</sup> *Suppose  $u$  is harmonic on an open square  $S$  and continuous on  $\bar{S}$ . Then its average on  $\partial S$  equals the average over the two diagonals.*

**PROOF.** By dilating to a subcube of  $S$  and then expanding back, we may assume that both  $u$  and its harmonic conjugate  $v$  are continuous on  $\bar{S}$ . Now  $S$  is composed of four  $15^\circ$  right triangles with common vertex the center of  $S$ . Let  $T$  be the lower triangle and denote its edges by  $L, B$  and  $R$ , where  $B$  is the hypotenuse. Applying Cauchy's theorem to  $u + iv$  on  $T$  and taking real parts of the integrals gives

$$0 = \oint_{\partial T} u dx - v dy = \int_B u - \frac{1}{\sqrt{2}} \left( \int_{R+L} u + \int_R v - \int_L v \right).$$

Using rotations and symmetry, applying this argument to the three remaining subtriangles of  $S$ , and summing the resulting equations, we see that the terms involving  $v$  cancel and we are left with our stated result.  $\square$

**THEOREM.** *If  $Nf \in L^1$ , then we may write  $f = \sum_j \lambda_j a_j$  so that the  $a_j$ 's are atoms and the coefficients  $\lambda_j$  satisfy*

$$(1) \quad \sum_j |\lambda_j| \leq 42 \|Nf\|_{L^1}.$$

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<sup>2</sup> This result appears to be a part of the folklore. In fact, Jim Davis has mentioned to the author that it appeared on his qualifying exam at Stanford. In the mean time, Joel Shapiro has been able to provide an alternate proof using elementary potential theory.

PROOF. Since  $u$  is continuous in  $\mathbf{R}_+^2$ , then  $E_k = \{x : Nf(x) > 2^k\}$  is open in  $\mathbf{R}$ . Let  $I(f) = \int_I f dx/|I|$  and define  $F_k$  as the complement in  $\mathbf{R}$  of  $E_k$ . We write  $E_k$  as the disjoint union of its collection  $\mathcal{C}_k$  of open components and then decompose  $f$  as a sum,  $f = g_k + h_k$ , where

$$(2) \quad g_k = \sum_{I \in \mathcal{C}_k} [f - I(f)]\chi_I, \quad h_k = f\chi_{F_k} + \sum_{I \in \mathcal{C}_k} I(f)\chi_I.$$

We claim that  $|h_k Z| \leq 7 \times 2^k$  a.e. Clearly, the estimate holds on  $F_k$  since  $Nf \leq 2^k$  there and  $|f| \leq Nf$  a.e. For the remaining set  $E_k$ , we fix an interval  $I$  in  $\mathcal{C}_k$  and show that

$$(3) \quad |I(f)| \leq 7 \times 2^k.$$

Let  $S_\varepsilon$  be the open square  $I \times (\varepsilon, |I| + \varepsilon)$  in  $\mathbf{R}_+^2$ . By the Proposition and letting  $\varepsilon \downarrow 0$ ,  $I(f)$  is seen to equal four times the average of  $u$  over the union of the two main diagonals less the sum of its averages over the three remaining sides. But the endpoints of  $I$  belong to  $F_k$ , so the diagonals, sides and top of  $S$  all belong to the “good” set for  $u$ , namely  $\Gamma = \{(z, y) \in \Gamma_x : x \in F_k\}$ . The definitions of  $F_k$  and  $Nf$  imply that  $u$  is bounded by  $2^k$  on  $\Gamma$  which establishes (3).

Following Coifman [1] and Herz [6f], the atoms are defined by

$$(4) \quad a_I \lambda_I^{-1} (g_k - g_{k+1})\chi_I, \quad \lambda_I = 21 \times 2^k |I|,$$

for each  $I \in \mathcal{C}_k$  and all integers  $k$ . By telescoping and using both that  $g_k - g_{k+1} = h_{k+1} - h_k$  and that  $g_{k+1}$  is supported in  $E_{k+1} \subset E_k$ , it follows that

$$f = \sum_{k=-\infty}^{\infty} (g_k - g_{k+1}) = \sum_k \sum_{I \in \mathcal{C}_k} \lambda_I a_I.$$

Each  $a_I$  is an atom since it is supported in  $I$  and the estimate  $\|a_I\|_\infty \leq |I|^{-1}$  follows from our  $L^\infty$  estimate on the  $h_k$ ,

$$\|g_k - g_{k+1}\|_\infty = \|h_{k+1} - h_k\|_\infty \leq 7(2^{k+1} + 2^k) = 21 \times 2^k.$$

To see that  $a_I$  has mean value zero, it suffices to write it in the form

$$a_I = \lambda_I^{-1} \left( [f - I(f)]\chi_I - \sum_{\substack{J \in \mathcal{C}_{k+1} \\ J \subset I}} [f - J(f)]\chi_J \right).$$

To establish inequality (1) (subject to relabeling) we use

$$(5) \quad \sum_k \sum_{I \in \mathcal{C}_k} |\lambda_I| + 21 \sum_k 2^k \sum_{I \in \mathcal{C}_k} |I| = 21 \sum_k 2^k |E_k| = 21 \sum_k (2^{k+1} - 2^k) |E_k|.$$

Indeed by (5), summation by parts, and the fact  $Nf > 2^k$  on  $E_k$ , we have

$$(6) \quad \sum_k \sum_{I \in \mathcal{C}_k} |\lambda_I| \leq 42 \sum_k 2^k |E_k \setminus E_{k+1}| \leq 42 \int Nf(x) dx. \quad \square$$

Fefferman, Rivière and Sagher [3] estimated the  $K$ -functional

$$K(f, t) = \inf\{\|g\|_{H^1} + t\|h\|_{L^\infty} : g \in H^1, h \in L^\infty, f = g + h\}$$

in terms of the “grand maximal” operator to describe interpolation spaces for the pair. We provide a description in terms of  $Nf$ . Let  $g^*$  denote the decreasing rearrangement of  $|g|$ .

COROLLARY (OF THE PROOF OF THE THEOREM). *If  $f$  belongs to  $H^1 + L^\infty$ , then*

$$(7) \quad K(f, t) \sim \int_0^t (Nf)^*(s) ds, \quad t > 0.$$

PROOF. The subadditivity of the integral operator in (7) implies that it is dominated by  $K(f, t)$ . To establish the opposite estimate, we fix  $t > 0$  and select an integer  $j$  so that  $2^{j-1} < (Nf)^*(t) \leq 2^j$ . From the constructions in (2), we see

$$(8) \quad g_j = \sum_{k=j}^{\infty} (g_k - g_{k+1}) = \sum_{k=j}^{\infty} \left( \sum_{I \in C_k} \lambda_I a_I \right).$$

The estimate  $\|h_j\|_\infty \leq 14(Nf)^*(t)$  follows by our selection of the index  $j$ , while

$$(9) \quad \|g_j\|_{H^1} \leq 42c \int_{E_j} Nf(x) dx \leq 42c \int_0^t (Nf)^*(s) ds$$

is derived as in (5)–(6) using the identity (8). Combining these estimates completes the proof.  $\square$

Minor modifications using  $p$ -atoms permit extension of these results to  $H^p$  spaces ( $\frac{1}{2} < p < 1$ ) on  $\mathbf{R}$ . Beginning with  $Nf$  and using classical techniques (theorems of Spanne-Stein and Hardy-Littlewood), these results provide a simplified approach to the various descriptions of  $H^p(\mathbf{R})$  (duality, grand maximal operator). By conformally mapping the unit disc onto  $\mathbf{R}_+^2$  and estimating the appropriate integrals obtained from the Proposition, one obtains the expected results for the circle. Exploiting a Fourier analytical technique of Calderón, Wilson has given a proof of the atomic decomposition into  $L^2$  atoms for higher dimensions in [8], while the condition  $N(u + iv) \in L^1$  is required in [7]. Finally, the author extends his thanks to Colin Bennett and Guido Weiss for valuable discussions related to this paper.

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